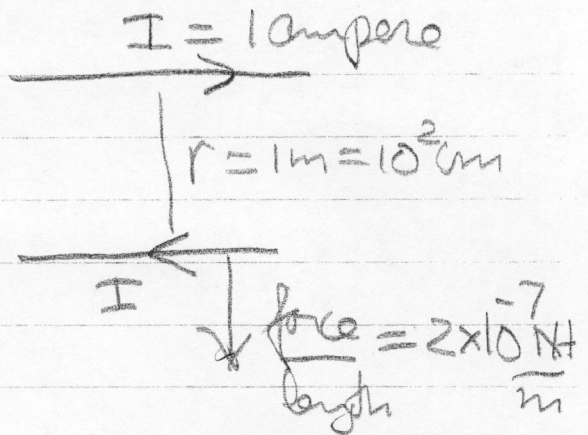


Solutions to homework problems

-1-

(1)



In Gaussian units,

$$\frac{\text{force}}{\text{length}} = \frac{2I^2}{c^2 r}$$

$$2 \times 10^{-7} \frac{(10^5 \text{ dyne})}{10^2 \text{ cm}} = \frac{2 (1 \text{ amp})^2}{9 \times 10^{20} \frac{\text{cm}^2}{\text{sec}^2} 10^2 \text{ cm}}$$

$$\therefore 9 \cdot 10^{-7+5-2+22} \frac{\text{dyne cm}^2}{\text{sec}^2} = (1 \text{ amp})^2$$

$\underbrace{\hspace{10em}}_{(\text{statamp})^2}$

$$3 \times 10^9 \text{ statamp} = 1 \text{ amp}$$

$$\longrightarrow \boxed{3 \times 10^9 \text{ stat coul} = 1 \text{ coul}}$$

$$\rightarrow \text{volt} = \frac{\text{Joule}}{\text{Coul.}} = \frac{10^7 \text{ erg}}{3 \times 10^9 \text{ statCoul}} = \frac{1}{3} 10^{-2} \text{ statvolts}$$

$$\text{henry} = \frac{\text{volt sec}}{\text{amp}} = \frac{1}{3} \frac{10^{-2} \text{ statvolt sec}}{3 \cdot 10^9 \text{ statamp}}$$

$$\rightarrow \text{henry} = \frac{1}{9} 10^{-11} \frac{\text{statvolt sec}}{\text{statamp}} = \frac{1}{9} 10^{-11} \frac{\text{sec}^2}{\text{cm}}$$

for $q = 1 \text{ coul}, B = 1 \text{ tesla}$

$$v = 1 \frac{\text{m}}{\text{sec}} \quad F = qvB = 1 \text{ NT}$$

$$F = q v c B$$

$$10^5 \text{ dyne} = \frac{3 \times 10^9 \text{ statcoul}}{3 \times 10^{10} \frac{\text{cm}}{\text{sec}}} 10^2 \frac{\text{cm}}{\text{sec}} (1 \text{ tesla})$$

$$\rightarrow 10^4 \frac{\text{dyne}}{\text{statcoul}} = 1 \text{ tesla} = \frac{\text{gauss}}{\text{cm}}$$

- 3 -

$$\text{Farad} = \frac{\text{Coul}}{\text{Volt}} = \frac{3 \times 10^9 \text{ statcoul}}{\frac{1}{3} 10^2 \text{ statvolt}}$$

$$= 9 \times 10^{11} \text{ cm}$$

2. Jackson 2.2

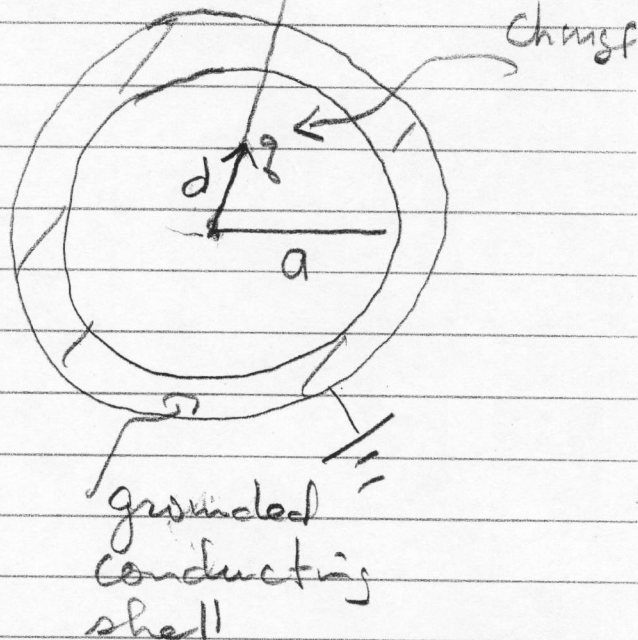
image $\rightarrow q' = -\frac{qa}{d}$

$D = a^2/d$

(a) for $r < a$

$$\phi(r) = \left\{ \frac{q}{|r-d|} - \frac{qa/d}{|r-D|} \right\}$$

$$D = \frac{a^2}{d}$$



(b) induced surface charge

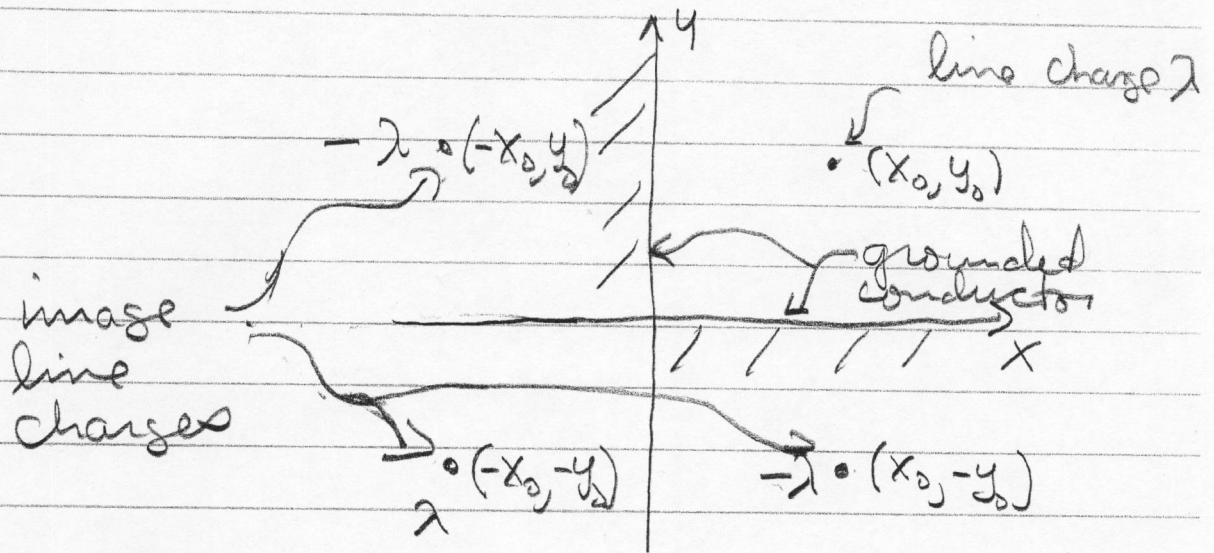
$$\sigma = -\frac{1}{4\pi} E_r(r=a) = +\frac{1}{4\pi} \left. \frac{\partial \phi}{\partial r} \right|_{r=a}$$

$$\sigma = \frac{-qq'}{|d-D|^2} = \frac{q^2 a/d}{|d - a^2/d|^2}$$

(d) $\phi \rightarrow \phi + V$, but \vec{E} and σ unchanged.

No changes if charge Q is added. Q goes to outside

3. Jackson 2.3



for $x \geq 0, y \geq 0$ (1st quadrant)

$$(a) \quad \phi = -\lambda \ln[(x-x_0)^2 + (y-y_0)^2] + \lambda \ln[(x-x_0)^2 + (y+y_0)^2] \\ - \lambda \ln[(x+x_0)^2 + (y+y_0)^2] + \lambda \ln[(x+x_0)^2 + (y-y_0)^2]$$

$$\phi = -\lambda \ln \left[\frac{[(x-x_0)^2 + (y-y_0)^2][(x+x_0)^2 + (y+y_0)^2]}{[(x-x_0)^2 + (y+y_0)^2][(x+x_0)^2 + (y-y_0)^2]} \right]$$

note that $\phi(x=0, y) = \phi(x, y=0) = 0$

$$\therefore \left. \frac{\partial \phi}{\partial y} \right|_{x=0} = \left. \frac{\partial \phi}{\partial x} \right|_{y=0} = 0$$

b.

$$\sigma = -\frac{1}{4\pi} \frac{\partial \phi}{\partial y} \Big|_{y=0} = +\frac{\lambda}{4\pi} \left\{ \frac{-2y_0}{(x-x_0)^2+y_0^2} \right.$$

$$\left. \frac{-2y_0}{(x-x_0)^2+y_0^2} + \frac{2y_0}{(x+x_0)^2+y_0^2} + \frac{2y_0}{(x+x_0)^2+y_0^2} \right\}$$

$$\frac{\sigma}{\lambda} = \frac{4y_0}{4\pi} \left\{ \frac{-1}{(x-x_0)^2+y_0^2} + \frac{1}{(x+x_0)^2+y_0^2} \right\}$$

$$\frac{\sigma}{\lambda} \frac{\pi x_0^2}{y_0} = \left\{ \frac{-1}{\left(\frac{x}{x_0}-1\right)^2 + \left(\frac{y_0}{x_0}\right)^2} + \frac{1}{\left(\frac{x}{x_0}+1\right)^2 + \left(\frac{y_0}{x_0}\right)^2} \right\}$$

only one parameter $\left(\frac{y_0}{x_0}\right)$ $\rightarrow g\left(\frac{x}{x_0}, \frac{y_0}{x_0}\right)$

note that $g \leq 0$ for $\frac{x}{x_0} > 0$

$g = 0$ for $\frac{x}{x_0} = 0, \infty$

c.

$$Q_x = \int_0^{\infty} dx \sigma_x = \frac{\lambda y}{\pi} \left\{ - \int_0^{\infty} \frac{dx}{(x-x_0)^2 + y_0^2} + \int_0^{\infty} \frac{dx}{(x+x_0)^2 + y_0^2} \right\}$$

$$= \frac{2y_0}{\pi} \left\{ - \int_{-x_0}^{\infty} \frac{dx}{x^2 + y_0^2} + \int_{x_0}^{\infty} \frac{dx}{x^2 + y_0^2} \right\}$$

$$= -\frac{2y_0}{\pi} \int_{x_0}^{x_0} \frac{dx}{x^2 + y_0^2} = -\frac{2\lambda y_0}{\pi} \frac{1}{y_0} \tan^{-1} \frac{x_0}{y_0}$$

d. For $\rho = \sqrt{x^2 + y^2} \gg \rho_0 = \sqrt{x_0^2 + y_0^2}$

expect quadrupole term in Taylor expansion for potential

$$\text{let } \lambda_1 = \lambda \quad (x_1, y_1) = (x_0, y_0)$$

$$\lambda_2 = -\lambda \quad (x_2, y_2) = (-x_0, y_0)$$

-13-

$$\lambda_3 = \lambda \quad (x_3, y_3) = (-x_0, -y_0)$$

$$\lambda_4 = -\lambda \quad (x_4, y_4) = (x_0, -y_0)$$

$$\Phi = - \sum_{i=1}^4 \lambda_i \ln \left(\frac{\rho - \rho_i}{\rho} \right)^2 \approx - \left(\sum_{i=1}^4 \lambda_i \right) \ln \rho^2$$

Taylor expand using $|\rho| \gg |\rho_i| \quad \parallel \quad 0$

$$- \left(\sum_{i=1}^4 \lambda_i \rho_i \right) \cdot \nabla_{\rho} \ln \rho^2$$

$$- \frac{1}{2} \left(\sum_{i=1}^4 x_i^2 \lambda_i \right) \frac{\partial^2 \ln \rho^2}{\partial x^2}$$

$$- \frac{1}{2} \left(\sum_{i=1}^4 y_i^2 \lambda_i \right) \frac{\partial^2 \ln \rho^2}{\partial y^2}$$

$$- \frac{1}{2} \left(\sum_{i=1}^4 2x_i y_i \lambda_i \right) \frac{\partial^2 \ln \rho^2}{\partial x \partial y}$$

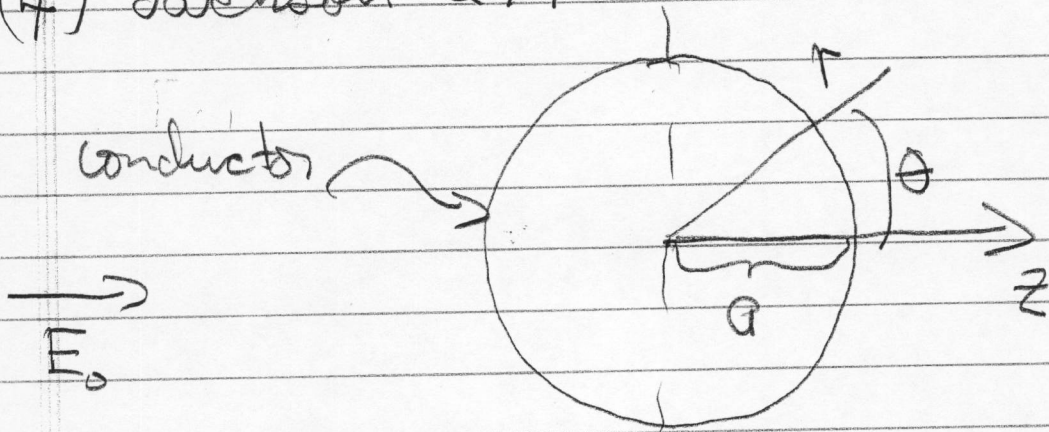
$\delta x_0 y_0 \lambda$

-14-

$$\varphi \approx -4x_0 y_0 \lambda \frac{\partial^2}{\partial x \partial y} \ln(x^2 + y^2)$$
$$\underbrace{\hspace{10em}}_{\frac{-4xy}{(x^2 + y^2)^2}}$$

$$\varphi \approx \frac{+16\lambda x_0 y_0 xy}{\rho^4}$$

(7) Jackson 2.9



(a) outside conductor

$$\varphi(r, \theta) = -E_0 r \cos \theta + \frac{E_0 a^3 \cos \theta}{r^2}$$

$$F_z = \int_0^{\pi/2} 2\pi a^2 \sin\theta d\theta \underbrace{\frac{1}{8\pi} \left(\frac{\partial\phi}{\partial r}\right)^2}_{\substack{\text{tension} \\ \text{in field} \\ \text{lines at} \\ \text{surface}}} \cos\theta$$

$$F_z = \int_0^{\pi/2} \frac{2\pi a^2 \sin\theta d\theta}{8\pi} \cos\theta \left[-E_0 \cos\theta - \frac{2E_0 a^3 \cos\theta}{a^3} \right]^2$$

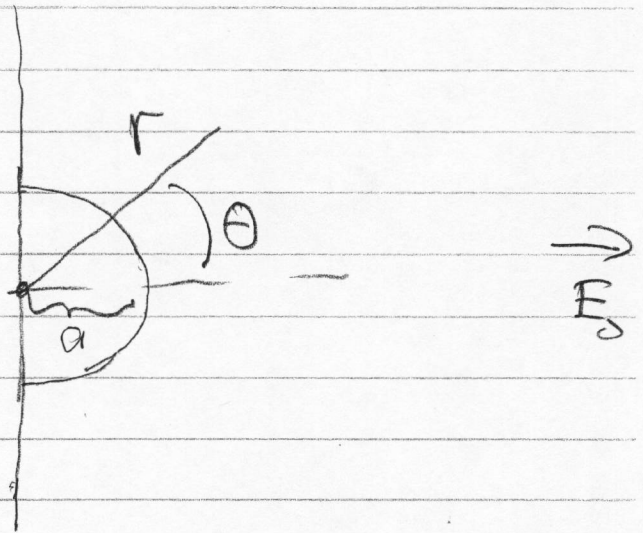
$$F_z = \frac{a^2}{4} 9E_0^2 \int_0^{\pi/2} \underbrace{\sin\theta d\theta}_{-d(\cos\theta)} \cos^3\theta = \frac{9}{16} a^2 E_0^2$$

(b) $\phi(r, \theta) = -E_0 r \cos\theta + \frac{E_0 a^3 \cos\theta}{r^2} + \frac{Q}{r}$

$$\left. \frac{\partial\phi}{\partial r} \right|_{r=a} = -3E_0 \cos\theta - \frac{Q}{a^2}, \text{ etc.}$$

(5) Jackson 2.10

(a)



Outside the hemispherical boss, the potential is given by

$$\phi(r, \theta) = -E_0 r \cos \theta + \frac{E_0 a^3 \cos \theta}{r^2}$$

on the boss

$$\sigma = -\frac{1}{4\pi} \frac{\partial \phi}{\partial r} = \frac{3E_0 \cos \theta}{4\pi}$$

on the plane

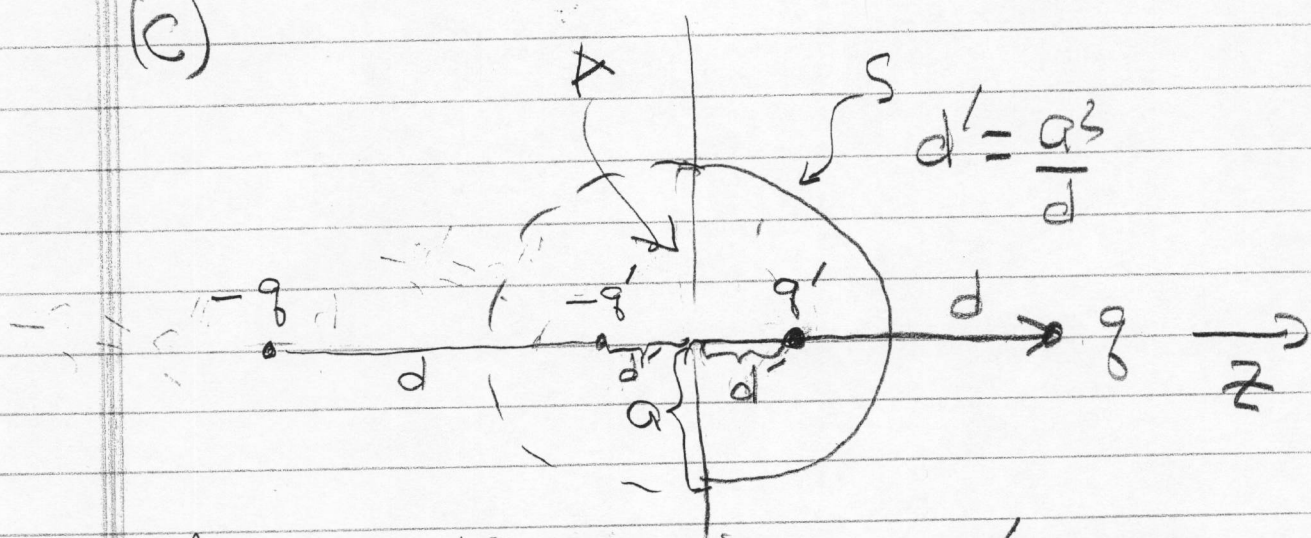
$$\sigma = -\frac{1}{4\pi} \frac{\partial \phi}{\partial z} \Big|_{z=0} = -\frac{1}{4\pi} \left(-\frac{\partial \phi}{\partial r} \right)_{\theta=\pi/2}$$

$$\sigma = \frac{1}{4\pi r} \left[E_0 r \sin\theta - \frac{E_0 a^3}{r^2} 3 \sin\theta \right] \quad \theta = \frac{\pi}{2}$$

$$\sigma = \frac{E_0}{4\pi} \left(1 - \frac{a^3}{r^3} \right) \quad \text{for } r > a$$

$$\begin{aligned} \text{(b)} \quad \Phi &= \int_0^{\pi/2} 2\pi a^2 \sin\theta d\theta \frac{3E_0 \cos\theta}{4\pi} \\ &= \frac{3}{2} E_0 a^2 \cdot \frac{1}{2} = \frac{3}{4} E_0 a^2 \end{aligned}$$

(c)



There are three images: $q' = -q \frac{a}{d}$, $-q'$, and $-q''$ (see figure)

The charge and three images make the potential zero on the plane and on the base.

$$\text{Charge on base} = q_b = \int_S \frac{\vec{E} \cdot d\vec{S}}{4\pi}$$

S is the hemispherical surface

Let A be the base of the hemisphere (a circular disc).

By Gauss' law

$$-q \frac{q}{d} = q' = \int_S \frac{\vec{E} \cdot d\vec{S}}{4\pi} + \int_A \frac{\vec{E} \cdot d\vec{S}}{4\pi}$$

$$\therefore q_b = -q \frac{q}{d} - \int_A \frac{\vec{E} \cdot d\vec{S}}{4\pi}$$

$$q_b = -q \frac{q}{d} - \int_0^a \frac{2\pi r dr}{4\pi} \left. \frac{\partial \phi}{\partial z} \right|_{z=0}$$

$$\frac{2q}{2z} = \frac{+2q d}{(d^2+r^2)^{3/2}} + \frac{2q' d'}{(d'^2+r^2)^{3/2}}$$

$$q_b = -q \frac{a}{d} - \frac{1}{2} \int_0^a r dr \left[\frac{2q d}{(d^2+r^2)^{3/2}} + \frac{2q' d'}{(d'^2+r^2)^{3/2}} \right]$$

$$q_b = -q \frac{a}{d} - \left[\frac{q d}{(d^2+r^2)^{1/2}} + \frac{q' d'}{(d'^2+r^2)^{1/2}} \right]_0^a$$

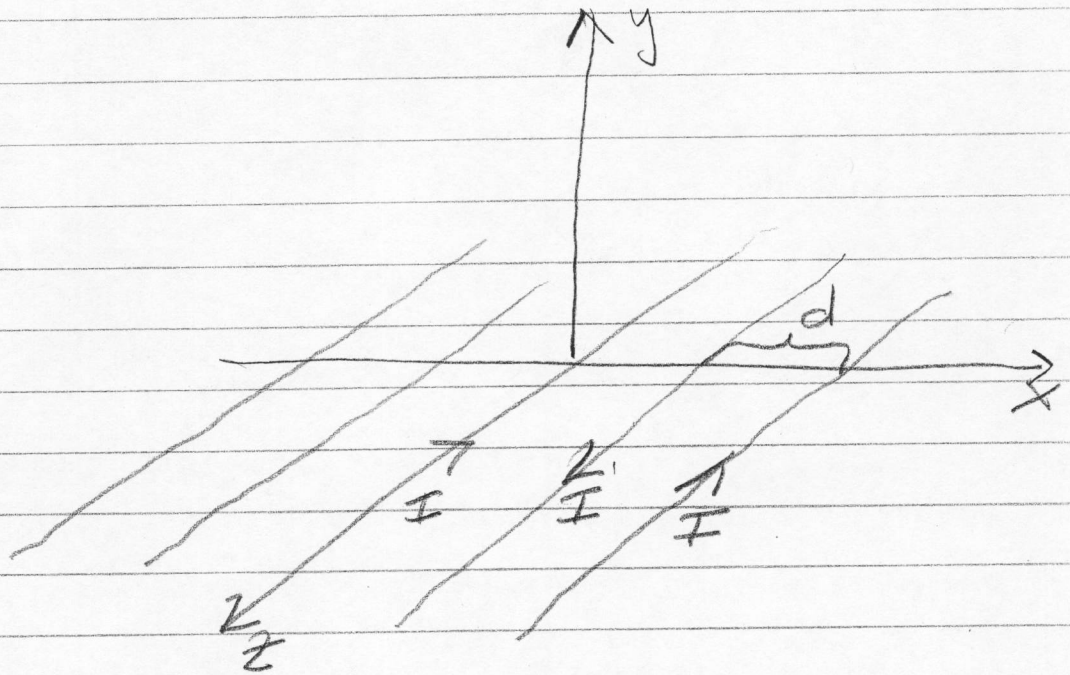
$$q_b = -q \frac{a}{d} - \left[1 + q' - \frac{2q d}{\sqrt{d^2+a^2}} - \frac{2q' d'}{\sqrt{d'^2+a^2}} \right]$$

$$q_b = -q \frac{a}{d} - q \left(1 - \frac{d}{\sqrt{d^2+a^2}} \right) + q' \frac{a}{d} \left(1 - \frac{a^2/d}{\sqrt{a^2/d^2+a^2}} \right)$$

$$q_b = -q + \frac{q}{\sqrt{d^2+a^2}} \left(+d - \frac{a^2}{d} \right)$$

$$q_b = -q + \frac{q}{\sqrt{d^2+a^2}} \left(\frac{d^2-a^2}{d} \right)$$

(6)



for $y > 0$ $\vec{j} = 0$ $\therefore \nabla \times \vec{B} = 0$

$\therefore \vec{B} = -\nabla \phi$

$0 = \nabla \cdot \vec{B} = -\nabla^2 \phi$

$\phi \rightarrow 0$ as $y \rightarrow \infty$

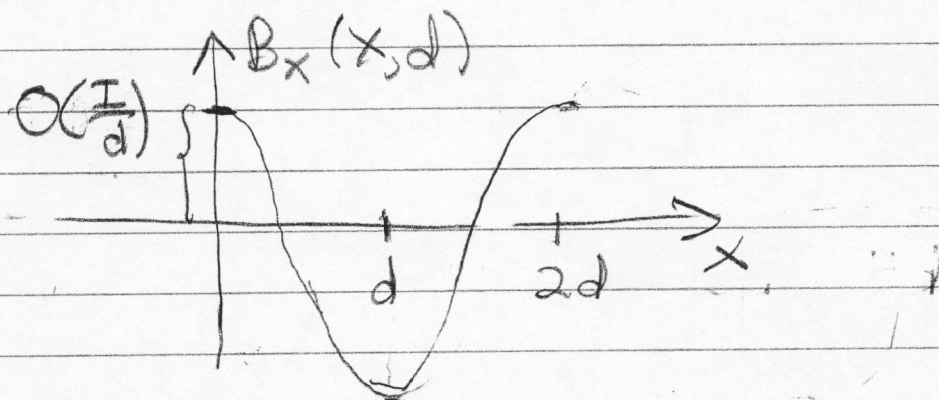
$\phi(x+2d, y) = \phi(x, y)$

$\therefore \phi(x, y) = \sum_n \left[A_n \cos\left(\frac{n\pi x}{d}\right) + B_n \sin\left(\frac{n\pi x}{d}\right) \right] e^{-\frac{n\pi y}{d}}$

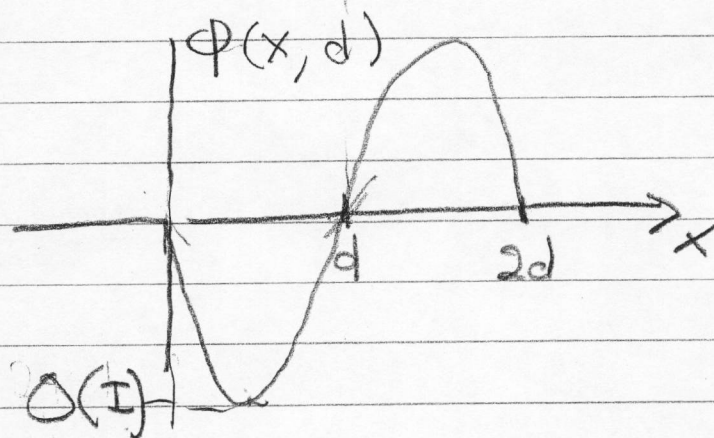
$$\int_0^{2d} \phi(x, y=d) \cos \frac{n\pi x}{d} dx = A_n \cdot 2d \cdot \frac{1}{2} e^{-\frac{\pi n d}{d}}$$

$$\int_0^{2d} \phi(x, y=d) \sin \frac{n\pi x}{d} dx = B_n \cdot d e^{-\pi n}$$

expect $B_x(x, d) = -\frac{\partial \phi(x, d)}{\partial x}$ to be of form



Hence



$$\therefore A_n = 0, \quad B_0 = 0, \quad B_1 \approx -I$$

for large y/d

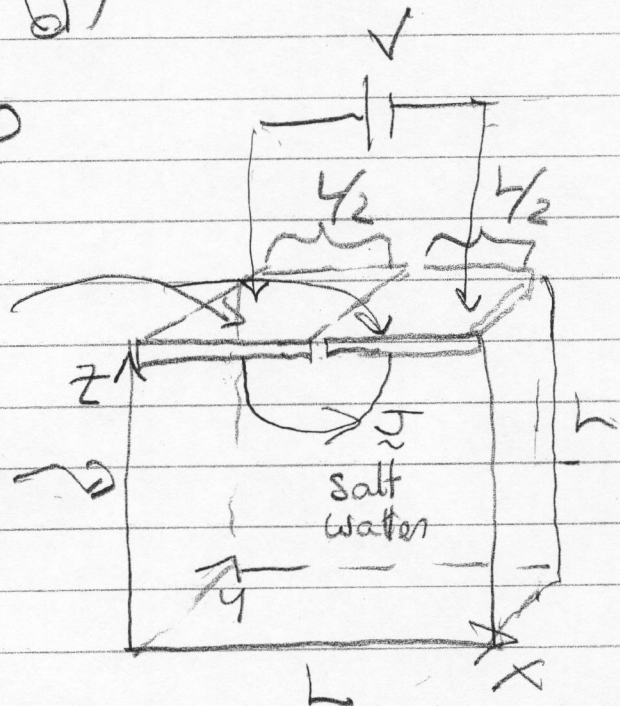
$$\Phi(x, y) \approx -I \sin\left(\frac{\pi y}{d}\right) e^{-\frac{\pi x}{d}}$$

$$\underline{B}(x, y) = -\underline{\nabla} \Phi$$

(7)

copper plates

cubic glass tank



in the salt water

$$\underline{E} = -\underline{\nabla} \Phi, \quad \underline{J} = -\sigma \underline{\nabla} \Phi$$

$$0 = \frac{\partial \rho}{\partial t} = -\underline{\nabla} \cdot \underline{J} = \sigma \nabla^2 \Phi$$

$$0 = \underline{J}_n = \sigma \frac{\partial \Phi}{\partial n} \quad \text{on glass walls}$$

$$\Phi(x, y, L) = \begin{cases} V & 0 < x < L/2 \\ 0 & L/2 < x < L \end{cases} \quad \left. \begin{array}{l} \text{on} \\ \text{lower} \\ \text{plates} \end{array} \right\}$$

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$$\Phi(x, y, z) = \sum_n A_n \cos\left(\frac{n\pi x}{L}\right) \cosh\left(\frac{n\pi z}{L}\right)$$

$$\left. \frac{\partial \Phi}{\partial x} \right|_{x=0} = \left. \frac{\partial \Phi}{\partial x} \right|_{x=L} = \left. \frac{\partial \Phi}{\partial y} \right|_{y=0} = \left. \frac{\partial \Phi}{\partial y} \right|_{y=L} = \left. \frac{\partial \Phi}{\partial z} \right|_{z=0} = 0$$

$$\int_0^L \Phi(x, y, z=L) \cos \frac{n'\pi x}{L} dx = A_{n'} \frac{L}{2} \cosh(\pi n')$$

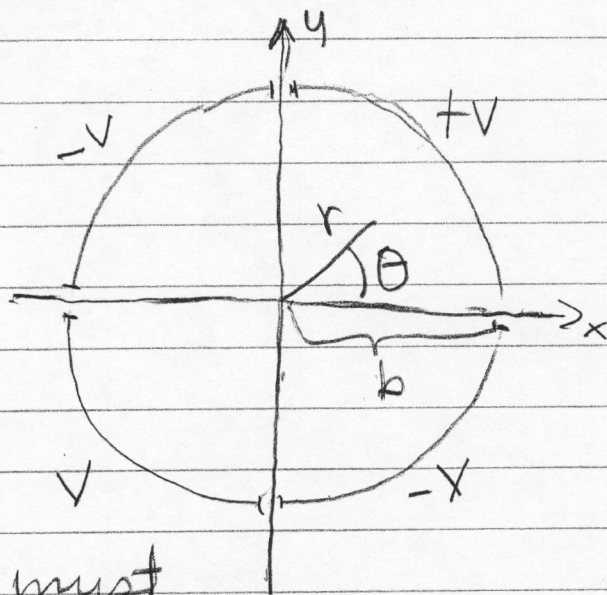
$$\sqrt{\int_0^{L/2} \cos\left(\frac{n'\pi x}{L}\right) dx} = A'_{n'} \frac{L}{2} \cosh(\pi n')$$

$$\sqrt{\frac{L}{n'\pi} \sin \frac{n'\pi x}{L} \Big|_0^{L/2}} = 11$$

$\underbrace{\hspace{10em}}_{\sin \frac{n'\pi}{2}}$

$$A_{n'} = \begin{cases} 0 & \text{for } n' \text{ even} \\ \frac{-\sqrt{2}(-)^{(n'+1)/2}}{n'\pi \cosh(\pi n')} & \text{for } n' \text{ odd} \end{cases}$$

(8) J. 2, 14 (a)



Inside cylinders, must discard solns varying like $\ln r$, $\frac{1}{r^l}$

Since $\phi(r, \theta)$ is odd in θ , need retain only $\sin l\theta$ terms

$$\phi(r, \theta) = \sum_{l=1}^{\infty} A_l r^l \sin(l\theta)$$

$$\int_0^{2\pi} d\theta \phi(b, \theta) \sin l\theta = A_l b^l 2\pi \frac{1}{2}$$

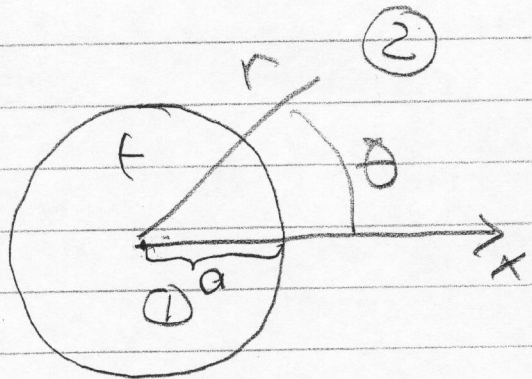
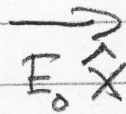
- 25 -

$$\begin{aligned}\pi b^l A_l &= 2V \int_0^{\pi/2} d\theta \sin \theta - 2V \int_{\pi/2}^{\pi} d\theta \sin \theta \\ &= \frac{2V}{l} \left[\left(1 - \cos \frac{2l\pi}{2} \right) + \cos l\pi - \cos \frac{l\pi}{2} \right] \\ &= \frac{2V}{l} \begin{cases} 0 & l \neq 2n+2 \\ 4 & l = 2n+2 \end{cases}\end{aligned}$$

where $n = 0, 1, 2, \dots$

$$\varphi(r, \theta) = \sum_{n=0}^{\infty} \frac{\pi b^l 2V}{(2n+2)} \left(\frac{r}{b} \right)^{2n+2} \sin(2n+2)\theta$$

(9)



$$\phi = \begin{cases} \phi_1(r, \theta) & \text{for } r < a \\ \phi_2(r, \theta) & \text{for } r > a \end{cases}$$

$$\phi_1(a, \theta) = \phi_2(a, \theta)$$

$$\neq \left. \frac{\partial \phi_1}{\partial r} \right|_{r=a} = \left. \frac{\partial \phi_2}{\partial r} \right|_{r=a}$$

$$\phi_2(r, \theta) \rightarrow -E_0 r \cos \theta \quad \text{for } r \rightarrow \infty$$

try

$$\phi_1(r, \theta) = A r \cos \theta$$

$$\phi_2(r, \theta) = -E_0 r \cos \theta + \frac{B}{r} \cos \theta$$

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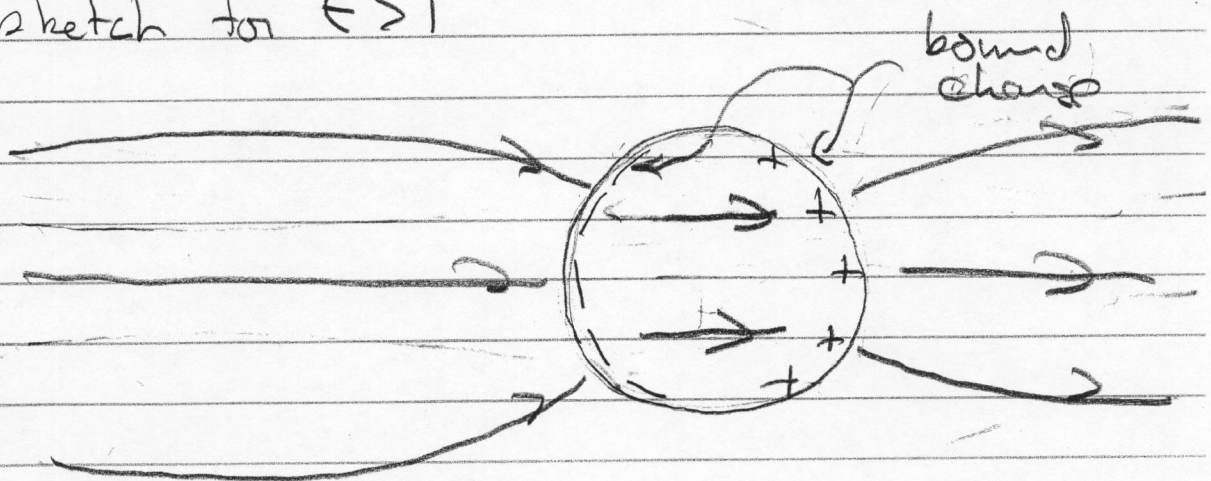
$$Aq = -E_0 a + \frac{B}{a^2}$$

$$\epsilon A = -E_0 - \frac{B}{a^2}$$

$$(\epsilon+1)A = -2E_0, \quad B = a^2 E_0 \left(\frac{-2}{1+\epsilon} + 1 \right) \\ = a^2 E_0 \left(\frac{\epsilon-1}{\epsilon+1} \right)$$

$$\therefore \varphi(r, \theta) = \begin{cases} \frac{-2E_0 r \cos\theta}{\epsilon+1} & r < a \\ -E_0 r \cos\theta + \frac{E_0 a^2 (\epsilon-1) \cos\theta}{r(\epsilon+1)} & r > a \end{cases}$$

sketch for $\epsilon > 1$

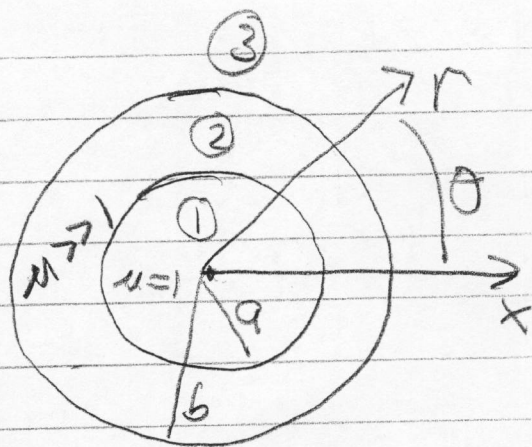


(10) $\nabla \times \vec{H} = 0$

$H = -\nabla \phi$

$\vec{B} = -\mu \nabla \phi$

$\mu = 1$



$$\phi(r, \theta) = \begin{cases} \phi_1 & \text{for } r < a \\ \phi_2 & \text{for } a < r < b \\ \phi_3 & \text{for } b < r \end{cases}$$

$\nabla^2 \phi_1 = \nabla^2 \phi_2 = \nabla^2 \phi_3 = 0$

$\phi_1(a, \theta) = \phi_2(a, \theta), \quad \phi_2(b, \theta) = \phi_3(b, \theta)$

$\vec{B} \cdot \hat{n}$ continuous across interface

b.c.

$\left. \frac{\partial \phi}{\partial r} \right|_a = \mu \left. \frac{\partial \phi_2}{\partial r} \right|_a, \quad \mu \left. \frac{\partial \phi_2}{\partial r} \right|_b = \left. \frac{\partial \phi_3}{\partial r} \right|_b$

as $r \rightarrow \infty \quad \phi_3(r, \theta) = -B_0 r \cos \theta$

$$\phi_1(r, \theta) = Ar \cos \theta$$

$$\phi_2(r, \theta) = Cr \cos \theta + \frac{D}{r} \cos \theta$$

$$\phi_3 = -B_0 r \cos \theta + \frac{E}{r} \cos \theta$$

$$Aa = Ca + \frac{D}{a^2}, \quad C + \frac{D}{b^2} = -B_0 + \frac{E}{b^2}$$

$$A = u \left(C - \frac{D}{a^2} \right), \quad u \left(C - \frac{D}{b^2} \right) = -B_0 - \frac{E}{b^2}$$

$$0 = C(u-1) - \frac{D}{a^2}(u+1), \quad C(u+1) + \frac{D(1-u)}{b^2} = -2B_0$$

$$D \left\{ \frac{(u+1)^2}{(u-1)a^2} + \frac{(1-u)}{b^2} \right\} = -2B_0$$

$$C = \frac{-2B_0(u+1)}{\left\{ \frac{(u+1)^2}{(u-1)} + \frac{a^2(1-u)}{b^2} \right\}(u-1)} = \frac{-2B_0(u+1)}{\left[(u+1)^2 - \frac{a^2(1-u)^2}{b^2} \right]}$$

$$A = \mu \left[C - \frac{D}{a^2} \right] = \mu \frac{-2B_0(\mu+1)}{\left[(\mu+1)^2 - \frac{a^2}{b^2}(\mu-1)^2 \right]}$$

$$\frac{2B_0(1-\mu)}{\left[(\mu+1)^2 - \frac{a^2}{b^2}(\mu-1)^2 \right]}$$

$$A = \frac{-2B_0 2\mu}{\left[(\mu+1)^2 - \frac{a^2}{b^2}(\mu-1)^2 \right]}$$

check limits

(1) $\mu = 1$

$A = -B_0$

$C = -B_0$

$D = 0$

$E = 0$

} correct

(2) $\mu \rightarrow \infty$ $a < b$

$A \rightarrow 0$

magnetic shielding