

Solution to Final Problem 8

FRW metric

$$(a) \quad ds^2 = -dt^2 + R(t)^2 \left[\frac{dr^2}{1-kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right]$$

$$dx = \frac{dr}{\sqrt{1-kr^2}} \Rightarrow$$

$$k=0$$

$$dx = dr \Rightarrow$$

$$\chi = r$$

$$\Rightarrow S(\chi) = r =$$

$$k=+1$$

$$dx = \frac{dr}{\sqrt{1-r^2}}$$

$$\Rightarrow \chi = \arcsin r$$

$$\Rightarrow r = \sin \chi$$

$$= S(\chi) = \sin \chi$$

$$k=-1$$

$$dx = \frac{dr}{\sqrt{1+r^2}}$$

$$\Rightarrow \chi = \operatorname{arcsinh} r$$

$$\Rightarrow r = S(\chi) = \sinh \chi$$

$$ds^2 = -dt^2 + R(t)^2 \left[d\chi^2 + S(\chi)^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

$$(b) \quad g_{tt} = -1, \quad g_{rr} = R(t)^2, \quad g_{\theta\theta} = R^2 S^2, \quad g_{\phi\phi} = R^2 S^2 \sin^2 \theta$$

(c) (a) Find conserved quantities & geodesic equations.

$$\text{Recall } \frac{du_\alpha}{d\tau} = \frac{1}{2} g_{\mu\nu,\alpha} u^\mu u^\nu$$

$$\text{where } u^\mu = dx^\mu/d\tau$$

$$\& \quad u_\alpha = g_{\alpha\beta} u^\beta$$

$$\text{For } \phi \text{ component: } g_{\mu\nu,\phi} = 0 \Rightarrow \frac{du_\phi}{d\tau} = 0$$

$$\Rightarrow u_\phi = \text{constant (conserved)}$$

$$u_\phi = g_{\phi\phi} u^\phi = R^2 S(\chi)^2 \sin^2 \theta u^\phi$$

$$\text{Now at } \chi=0, \quad S(\chi)=0 \quad \therefore \quad u_\phi=0 \quad \text{at } \chi=0$$

Since u_ϕ is constant along the geodesic $u_\phi=0$ everywhere.

$$\text{If } u_\phi=0 \Rightarrow u^\phi=0 \Rightarrow u^\phi = \frac{d\phi}{d\tau} = 0 \Rightarrow \phi = \text{constant along the geodesic}$$

For θ component: only $g_{\theta\theta}$ depends on θ

$$\frac{du_\theta}{d\tau} = \frac{1}{2} g_{\theta\theta,\theta} u^\theta u^\theta = 2 \sin \theta \cos \theta (u^\theta)^2 = 0 \quad \text{since } u^\theta=0$$

$$u_\theta = \text{constant along geodesic}$$

$$u_\theta = g_{\theta\theta} u^\theta = R^2 S^2 u^\theta \Rightarrow u_\theta = 0 \text{ at } \chi=0 \Rightarrow u_\theta = 0 \text{ along geodesic (conserved)}$$

$$u_\theta = 0 \Rightarrow u^\theta = 0 \Rightarrow \frac{d\theta}{d\tau} = 0 \Rightarrow \theta = \text{constant along geodesic}$$

χ component

$$\frac{du_\chi}{d\tau} = \frac{1}{2} g_{\mu\nu, \chi} u^\mu u^\nu$$

We know from before that $u^\theta = u^\phi = 0$

$$= \frac{1}{2} g_{\theta\theta, \chi} u^\theta u^\theta + \frac{1}{2} g_{\phi\phi, \chi} u^\phi u^\phi + \frac{1}{2} g_{\theta\theta, \chi} u^\theta u^\theta + \frac{1}{2} g_{\phi\phi, \chi} u^\phi u^\phi$$

$$= \frac{1}{2} R^2 \frac{d^2 \theta}{d\chi} u^\theta u^\theta + \frac{1}{2} R^2 \sin^2 \theta \frac{d^2 \phi}{d\chi} u^\phi u^\phi$$

$$= 0 \quad \text{since } u^\theta = u^\phi = 0$$

$\Rightarrow \frac{du_\chi}{d\tau} = 0 \Rightarrow u_\chi = \text{constant along geodesic}$

$$u_\chi = g_{\chi\chi} u^\chi = \left[R^2 \frac{d\chi}{d\tau} = \text{constant} \right]$$

$$\Rightarrow \frac{d\chi}{d\tau} = \frac{\text{constant}}{R^2}$$

t component

$u^\chi u_\chi = -1$ normalization

$$-u^t u_t + u^r u_r + u^\theta u_\theta + u^\phi u_\phi = -1$$

$$-\left(\frac{dt}{d\tau}\right)^2 + g_{rr} u^r u^r = -1$$

u

geodesic equation

$$-\left(\frac{dt}{d\tau}\right)^2 + \frac{(R^2 \frac{d\chi}{d\tau})^2}{R^2} = -1 \Rightarrow \left(\frac{dt}{d\tau}\right)^2 = 1 + R^2 \left(\frac{d\chi}{d\tau}\right)^2$$

② Interpret $\frac{d\chi}{d\tau} = \frac{\text{const}}{R^2}$ for $k=0$ case.

$$\text{For } k=0, \chi=r \Rightarrow \frac{dr}{dt} = \frac{\text{const}}{R(t)^2}$$

If particle starts with peculiar velocity $v_p = \frac{dr}{dt}$ that velocity will reduce as $R(t)$ expands.