

PHYS 2D
DISCUSSION SECTION

2012/5/16

- 
- Quiz 3 is graded
 - Pick up quiz 3 today, tomorrow, or next Tuesday
 - Regrade request: 1 week

Topics



- Probability
- Born interpretation
- Normalization
- Operators
- x & p operators
- Schrodinger's equation
- Free space wave function
- Particle in a box & Finite square well
- Quantum oscillator

Probability

- 2 possible outcomes for a measurement:
- ◆ Event 1 has probability p_1 of happening, event 2 has p_2 . When taking a measurement, either 1 happens or 2 happens, so $p_1 + p_2 = 1$
- ◆ 1 measurement: has a chance p_1 to find result 1, chance p_2 to find result 2
- ◆ 100 measurements: find result 1 $100 * p_1$ times, result 2 $100 * p_2$ times
- ◆ N measurements: result in 1 Np_1 times, 2 Np_2 times
- ◆ $Np_1 + Np_2 = N(p_1 + p_2) = N$

Probability



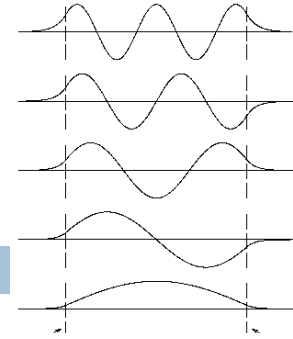
- M possible outcomes for a measurement:
- ◆ A measurement has chance p_x to return result x , where x goes from 1 to M , and $\sum_x p_x = 1$
- If M is infinite:
- ◆ Consider for example $x = \text{position of a particle}$
- ◆ x is now continuous (infinite values of x)
- ◆ If we assign each x a probability, sum over x blows up
- ◆ Define $p(x)$ so that a measurement has chance $p(x)dx$ of finding the particle in an interval dx about x
- ◆ $p(x)$ is now a probability distribution

Probability



- $P(x)dx$ is the probability of finding the particle in a small interval dx around x
- If at each x a function f has value $f(x)$, then the average measured value of f after lots of measurements is $\bar{f} = \int f P(x) dx$
- Discrete example: Dice
 - ◆ $x=1$ to 6
 - ◆ $p_x=1/6$
 - ◆ $f_x=x$
 - ◆ Average measured value of $x = \sum_x f_x p_x = 3.5$

Wave function Ψ



- Wave function $\Psi(x)$ is used to describe a particle
- ◆ It contains all the information about that particle
- ◆ It is a complete description
- ◆ In principle, if we know $\Psi(x)$, we can deduce any physical quantity of the particle we want to know

Born interpretation

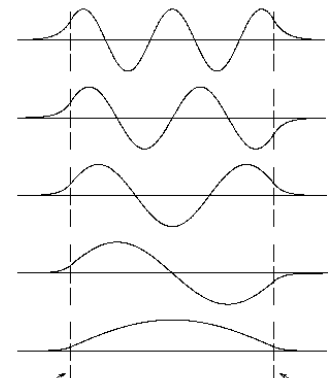


- Born interpretation: (the physical meaning of the wave function $\Psi(x)$)
- ◆ $|\Psi(x)|^2$ is a probability distribution, which means a measurement (of position x) has a chance $|\Psi(x)|^2 dx$ of returning a value in the interval dx about x ,

or

$$P(x)dx = |\Psi(x,t)|^2 dx$$

- ◆ When measuring the position of the particle, the returned value x , which is the position of the particle, can be anywhere from $-\infty$ to ∞ as long as $|\Psi(x)|^2 \neq 0$, but the probability for each x is different



Normalization



- If our $\Psi(x)$ describes 1 particle, then the probability of finding the particle at each x must add up to 1
- For discrete x , $\sum_x p_x = 1$
- For continuous x , $\int_{-\infty}^{\infty} |\Psi(x,t)|^2 dx = 1$ ($P(x)dx = |\Psi(x,t)|^2 dx$)
- This is the normalization condition, arising from the probabilistic nature of wave functions

Operators



- Wave function contains all physical quantities
 - ◆ Position, momentum, energy, etc.
- To extract these information (observables), need to define operators for the corresponding measurable physical quantity f
- Expectation value \bar{f} = average value of f after large number of measurements (of f)
- The operator for f , \hat{f} is defined so that the expectation value is $\bar{f} = \int_{-\infty}^{\infty} \psi^*(\mathbf{x}) \hat{f} \psi(\mathbf{x}) d\mathbf{x}$
- \hat{f} can be a number or a differential operator

x & p operators

□ For example, if f is the position x

◆ By definition, $\bar{x} = \int_{-\infty}^{\infty} \Psi^*(x) \hat{x} \Psi(x) dx$

◆ But $\bar{x} = \int_{-\infty}^{\infty} x P(x) dx = \int_{-\infty}^{\infty} x |\Psi(x)|^2 dx$
 $= \int_{-\infty}^{\infty} x \Psi^*(x) \Psi(x) dx = \int_{-\infty}^{\infty} \Psi^*(x) x \Psi(x) dx$

◆ So $\hat{x} = x$

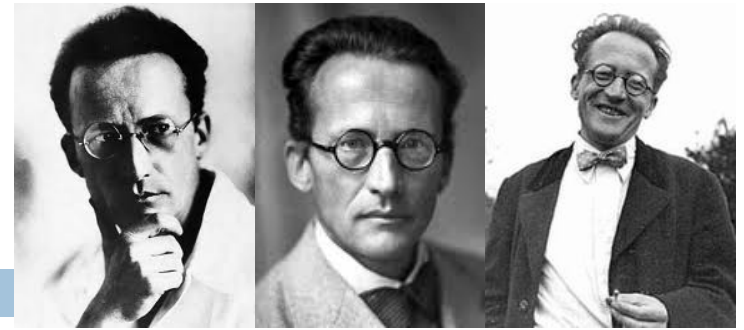
□ For momentum p , \hat{p} is a differential operator

◆ It can be derived: $\hat{p} = \frac{\hbar}{i} \frac{d}{dx}$

◆ So $\bar{p} = \int_{-\infty}^{\infty} \Psi^*(x) \hat{p} \Psi(x) dx = \int_{-\infty}^{\infty} \Psi^*(x) \frac{\hbar}{i} \frac{d}{dx} \Psi(x) dx$

◆ $\hat{p} = \frac{\hbar}{i} \frac{d}{dx}$ acts on the wave function $\Psi(x)$ to the right

Schrodinger's equation



- Relates the wave function of a particle to the environment ($U(x)$) of the particle

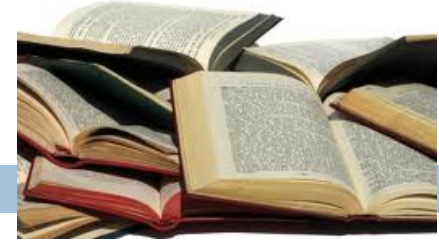
$$i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{x}, t) = \hat{H} \Psi(\mathbf{x}, t)$$

- The equation that determines the wave function
- Total energy $E = p^2/2m + U(x)$
- Convert p to \hat{p} and x to \hat{x} to get \hat{H}

- $$\hat{H} = \frac{\hat{p}^2}{2m} + U(\hat{x}) = \frac{\left(\frac{\hbar}{i} \frac{d}{dx}\right)^2}{2m} + U(\hat{x}) = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + U(\hat{x})$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x,t)}{\partial x^2} + U(x)\Psi(x,t) = i\hbar \frac{\partial \Psi(x,t)}{\partial t}$$

Schrodinger's equation



- Time-dependent Schrodinger's equation:

$$i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{x}, t) = \hat{H} \Psi(\mathbf{x}, t)$$

where \hat{H} is the total energy operator (Hamiltonian)

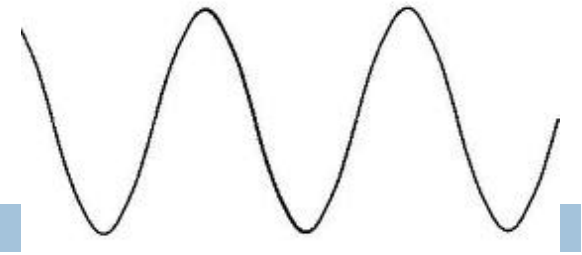
- ◆ If we assume that $\Psi(\mathbf{x}, t)$ can be separated

$$\Psi(\mathbf{x}, t) = \psi(\mathbf{x}) \phi(t)$$

- ◆ Then we have the time-independent Schrodinger's equation $\hat{H} \psi(\mathbf{x}) = E \psi(\mathbf{x})$, where E is the total energy

- ◆ Also $i\hbar \frac{\partial}{\partial t} \phi(t) = E \phi(t)$

Free space wave function



- Putting Schrodinger's equation to use
- Free space: $U(x)=0$

$$\hat{H} \psi (x) = E \psi (x)$$

$$i\hbar \frac{\partial}{\partial t} \phi (t) = E \phi (t)$$

$$\hat{H} = \frac{\hat{p}^2}{2m} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

$$\frac{d^2}{dx^2} \cos (x) = -\cos (x)$$

$$\frac{d^2}{dx^2} \sin (x) = -\sin (x)$$

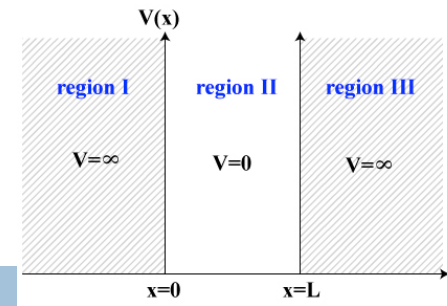
$$e^{ix} = \cos (x) + i \sin (x)$$

□ So, $\psi (x) = A e^{\pm i k x}$, $E = \frac{\hbar^2 k^2}{2m}$

$$\phi (t) = B e^{-i \omega t}$$
, $E = \hbar \omega$

$$\Psi (x, t) = C e^{i (\pm k x - \omega t)} = C [\cos (\pm k x - \omega t) + i \sin (\pm k x - \omega t)]$$

Particle in a box



- Not so trivial application of Schrodinger's equation
- $\phi(t) = Be^{-i\omega t}$ doesn't affect observables
- 1-D box, particle is restricted so $x \in [0, L]$
- To prevent particle from moving beyond $[0, L]$, we let $V = \infty$ outside the box
- When $V = \infty$, $\psi(x) = 0$, the wall is infinitely high
- To find $\psi(x)$ in region II, consider the time-independent Schrodinger's equation and require $\psi(x)$ be continuous ($\psi(0) = \psi(L) = 0$)

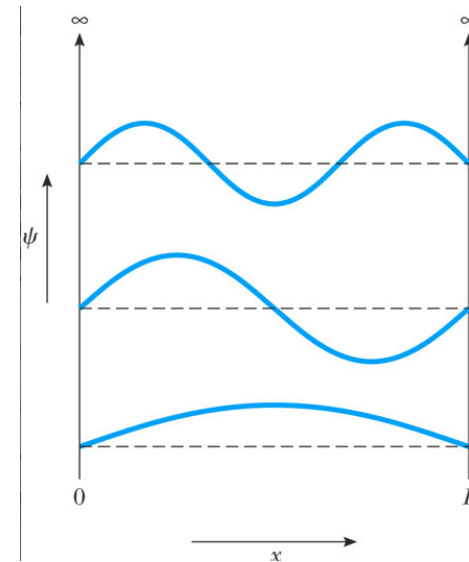
Particle in a box

- $\psi(x) = A \sin kx = A \sin[(n\pi/L)x]$

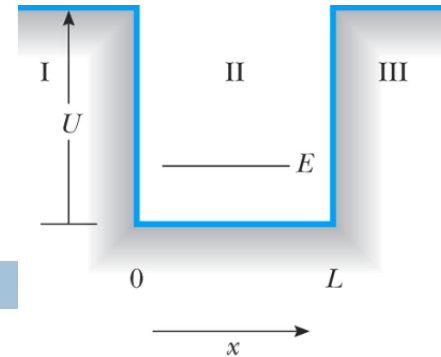
- $E = \frac{n^2 \pi^2 \hbar^2}{2mL}$

- $n = 1, 2, 3, \dots$

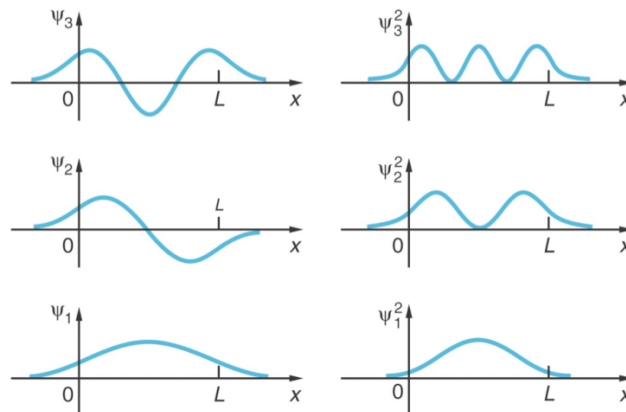
- Integer n (quantization) comes from boundary conditions



Finite square well

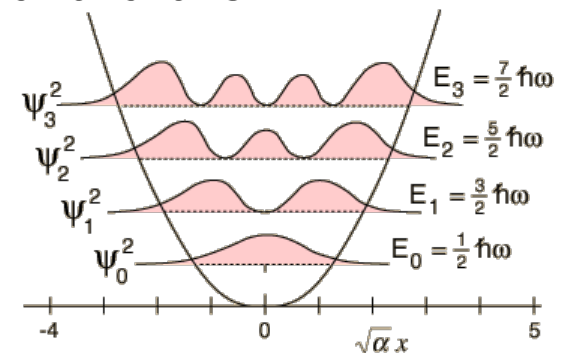


- The potential outside II is now $V=U$, not ∞
- Wave function similar to particle in a box
- Significant difference at the boundary ($x=0$ or L)
- Wave function is non-zero outside II, for a small distance (penetration depth)
- To find $\psi(x)$, require ψ and $d\psi/dx$ be continuous at $x=0$ & $x=L$



Quantum oscillator

- Classical oscillator: $E = p^2/2m + U(x)$, $U(x) = m\omega^2 x^2/2$
- Ex. Spring
- Quantum version: $-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + \frac{m\omega^2 x^2}{2} \psi(x) = E\psi(x)$
- $E = (n + 1/2)\hbar\omega$, $n = 0, 1, 2, \dots$
- $n = 0$ correspond to zero classical amplitude, but $E = \hbar\omega/2 \neq 0$
- Can model small oscillations around stable equilibrium points





□ Questions?