

PHYSICS 152B/232
Spring 2017
Homework Assignment #1 Solutions

[1] Consider a one-dimensional chain of s -orbitals

$$H = \sum_n \left(\varepsilon_A |A_n\rangle\langle A_n| + \varepsilon_B |B_n\rangle\langle B_n| \right. \\ \left. - t \sum_n \left(|A_n\rangle\langle B_n| + |B_n\rangle\langle A_{n+1}| + |B_n\rangle\langle A_n| + |A_{n+1}\rangle\langle B_n| \right) \right) .$$

- (a) How many atoms are there per unit cell? What is the length of the Wigner-Seitz cell?
- (b) Find the dispersions $E_a(k)$ of the energy bands.
- (c) Sketch the band structure over the one-dimensional Brillouin zone.
- (d) Show that for $\varepsilon_A = \varepsilon_B$ that you recover the correct energy band for the uniform one-dimensional nearest-neighbor chain.

Solution :

(a) There are two atoms per unit cell (one A and one B). The length of the Wigner-Seitz cell is $a = 2a_0$, where a_0 is the separation between neighboring A and B sites.

(b) From the Hamiltonian above, we read off the hopping matrix

$$H_{aa'}(n - n') = \begin{pmatrix} -\varepsilon_A \delta_{n-n',0} & t (\delta_{n-n',0} + \delta_{n-n',1}) \\ t (\delta_{n-n',0} + \delta_{n-n',-1}) & -\varepsilon_B \delta_{n-n',0} \end{pmatrix}$$

Thus,

$$\hat{H}_{aa'}(k) = \sum_j t_{aa'}(j) e^{-ikja} = \begin{pmatrix} \varepsilon_A & -t(1 + e^{-ika}) \\ -t(1 + e^{ika}) & \varepsilon_B \end{pmatrix}$$

The energy eigenvalues are then

$$E_{\pm}(k) = \frac{1}{2}(\varepsilon_A + \varepsilon_B) \pm \sqrt{\frac{1}{4}(\varepsilon_A - \varepsilon_B)^2 + 4t^2 \cos^2\left(\frac{1}{2}ka\right)} ,$$

where we've used $|1 + e^{-ika}|^2 = 2 + 2 \cos(ka) = 4 \cos^2\left(\frac{1}{2}ka\right)$.

(c) See the plots in Fig. 1.

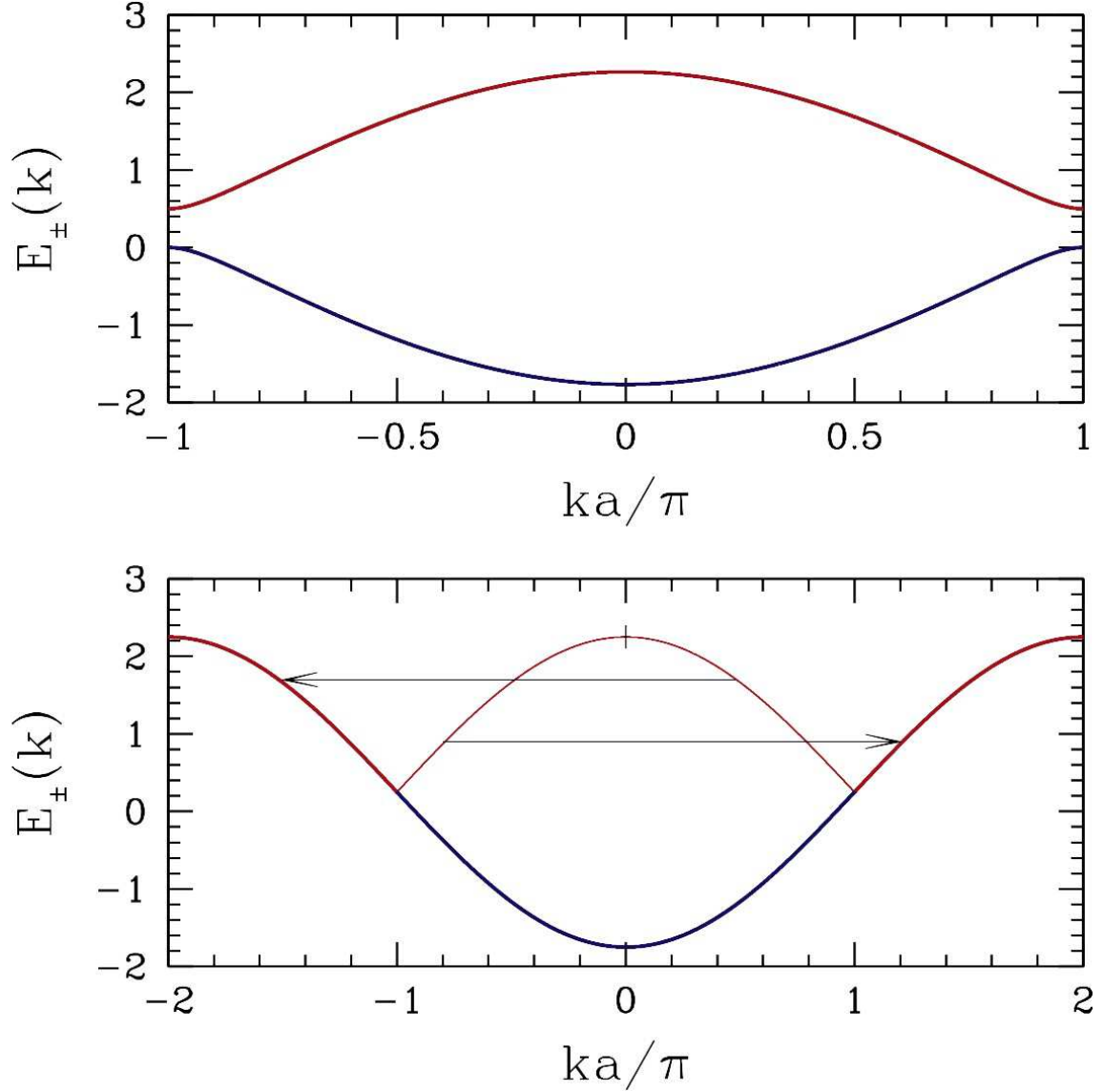


Figure 1: Energy bands $E_{\pm}(k)$ for problem 1. Top: $\varepsilon_A = 0.5$, $\varepsilon_B = 0.0$, $t = 1.0$. Bottom: $\varepsilon_A = \varepsilon_B = 0.5$, $t = 1$. The extended zone plot of the dispersion is shown for the latter case, in which the one-dimensional dispersion $E(k) = \varepsilon_0 - 2t \cos(ka')$ is recovered, with $a' = \frac{1}{2}a$ and $k \in [-\frac{\pi}{a'}, \frac{\pi}{a'}]$.

(d) Let $\varepsilon_A = \varepsilon_B = \varepsilon_0$. Then $E_{\pm}(k) = \varepsilon_0 \pm 2t \cos(\frac{1}{2}ka)$. If we translate the section of the + band on the interval $k \in [-\frac{\pi}{a}, 0]$ by $\frac{2\pi}{a}$, and the section on the interval $k \in [0, \frac{\pi}{a}]$ by $-\frac{2\pi}{a}$, we obtain the dispersion $E(k) = \varepsilon_0 - 2t \cos(ka')$ on the interval $k \in [-\frac{\pi}{a'}, \frac{\pi}{a'}]$, with $a' = \frac{1}{2}a$.

[2] Hexagonal boron nitride, BN, has a honeycomb lattice structure, with boron atoms at A sites and nitrogen atoms at B sites. The tight binding Hamiltonian is

$$H = \sum_{\mathbf{R}} \left(\varepsilon_{\text{A}} |A_{\mathbf{R}}\rangle \langle A_{\mathbf{R}}| + \varepsilon_{\text{B}} |B_{\mathbf{R}}\rangle \langle B_{\mathbf{R}}| \right) - t \sum_{\mathbf{R}} \left(|A_{\mathbf{R}}\rangle \langle B_{\mathbf{R}}| + |A_{\mathbf{R}}\rangle \langle B_{\mathbf{R}+\mathbf{a}_1}| + |A_{\mathbf{R}}\rangle \langle B_{\mathbf{R}-\mathbf{a}_2}| + \text{H.c.} \right) .$$

- (a) Find the 2×2 Hamiltonian matrix $\hat{H}(\mathbf{k})$. You may find it convenient to write $\mathbf{k} = \frac{\theta_1}{2\pi} \mathbf{b}_1 + \frac{\theta_2}{2\pi} \mathbf{b}_2$ and express your answer in terms of $\theta_{1,2}$.
- (b) Find expressions for the band energies at high symmetry points Γ , K, and M.
- (c) Find an expression for the band gap Δ . Is the gap direct or indirect?

Solution :

(a) The Hamiltonian matrix is

$$\hat{H}(\mathbf{k}) = \begin{pmatrix} \varepsilon_{\text{A}} & -t\gamma(\mathbf{k}) \\ -t\gamma^*(\mathbf{k}) & \varepsilon_{\text{B}} \end{pmatrix} ,$$

where

$$\gamma(\mathbf{k}) = 1 + e^{i\mathbf{k}\cdot\mathbf{a}_1} + e^{-i\mathbf{k}\cdot\mathbf{a}_2} = 1 + e^{i\theta_1} + e^{-i\theta_2} .$$

There are two bands:

$$E_{\pm}(\mathbf{k}) = \frac{1}{2}(\varepsilon_{\text{A}} + \varepsilon_{\text{B}}) \pm \sqrt{\frac{1}{4}(\varepsilon_{\text{A}} - \varepsilon_{\text{B}})^2 + t^2|\gamma(\mathbf{k})|^2} .$$

(b) Recall $\mathbf{k}_{\Gamma} = 0$, $\mathbf{k}_{\text{K}} = \frac{1}{3}\mathbf{B}_1 + \frac{1}{3}\mathbf{B}_2$, and $\mathbf{k}_{\text{M}} = \frac{1}{2}\mathbf{b}_1$. Thus,

$$\gamma(\Gamma) = 3 \quad , \quad \gamma(\text{K}) = 0 \quad , \quad \gamma(\text{M}) = 1$$

and

$$\begin{aligned} E_{\pm}(\Gamma) &= \frac{1}{2}(\varepsilon_{\text{A}} + \varepsilon_{\text{B}}) \pm \sqrt{\frac{1}{4}(\varepsilon_{\text{A}} - \varepsilon_{\text{B}})^2 + 9t^2} \\ E_{\pm}(\text{K}) &= \varepsilon_{\text{A}} , \varepsilon_{\text{B}} \\ E_{\pm}(\text{M}) &= \frac{1}{2}(\varepsilon_{\text{A}} + \varepsilon_{\text{B}}) \pm \sqrt{\frac{1}{4}(\varepsilon_{\text{A}} - \varepsilon_{\text{B}})^2 + t^2} . \end{aligned}$$

(c) Since nitrogen has a greater nuclear charge, we expect $\varepsilon_B < \varepsilon_A$. The maximum valence (–) band energy is then ε_B , at K. The minimum conduction band energy is ε_A , also at K. Thus, the gap is direct and equal to $\Delta = \varepsilon_A - \varepsilon_B$. The direct gap at wavevector \mathbf{k} is

$$\Delta(\mathbf{k}) \equiv E_+(\mathbf{k}) - E_-(\mathbf{k}) = \sqrt{(\varepsilon_A - \varepsilon_B)^2 + 4t^2|\gamma(\mathbf{k})|^2} \quad .$$

[3] Consider a tight binding model of (p_x, p_y) orbitals on a triangular lattice. The hopping is restricted to nearest neighbor links. Recall that the hopping matrix elements are given by

$$t_{\mu\nu}(\hat{\boldsymbol{\eta}}) = t_w \delta_{\mu\nu} - (t_s + t_w) \hat{\eta}_\mu \hat{\eta}_\nu \quad ,$$

where the link direction is $\hat{\boldsymbol{\eta}}$.

- (a) Find the matrix $\hat{t}_{\mu\nu}(\mathbf{k})$. You may find it convenient to write $\mathbf{k} = \frac{\theta_1}{2\pi} \mathbf{b}_1 + \frac{\theta_2}{2\pi} \mathbf{b}_2$ and express your answer in terms of $\theta_{1,2}$.
- (b) Find expressions for the band energies at the high symmetry points Γ , K, and M.
- (c) For $t_s = 1$ and $t_w = \frac{1}{2}$, plot the dispersions $E_\pm(\mathbf{k})$ along the path $\Gamma\text{MK}\Gamma$.

Solution :

$\hat{\boldsymbol{\eta}}$	$\hat{\eta}_j^x \hat{\eta}_j^x$	$\hat{\eta}_j^x \hat{\eta}_j^y$	$\hat{\eta}_j^y \hat{\eta}_j^y$
$\pm \hat{\mathbf{a}}_1$	$\frac{1}{4}$	$-\frac{\sqrt{3}}{4}$	$\frac{3}{4}$
$\pm \hat{\mathbf{a}}_2$	$\frac{1}{4}$	$\frac{\sqrt{3}}{4}$	$\frac{3}{4}$
$\pm \hat{\mathbf{a}}_3$	1	0	0

Table 1: Values of $\eta_j^\mu \eta_j^\nu$ for the six nearest neighbor vectors.

(a) On the triangular lattice, there are six nearest neighbors. Defining the primitive direct lattice vectors $\mathbf{a}_1 = a(\frac{1}{2}\hat{\mathbf{x}} - \frac{\sqrt{3}}{2}\hat{\mathbf{y}})$ and $\mathbf{a}_2 = a(\frac{1}{2}\hat{\mathbf{x}} + \frac{\sqrt{3}}{2}\hat{\mathbf{y}})$, the six nearest neighbor vectors are $\pm\mathbf{a}_1$, $\pm\mathbf{a}_2$, and $\pm\mathbf{a}_3$, where $\mathbf{a}_3 \equiv -\mathbf{a}_1 - \mathbf{a}_2 = -a\hat{\mathbf{x}}$. From the entries in Tab. 1, we have

$$\begin{aligned} \hat{t}_{xx}(\mathbf{k}) &= 2t_w(c_1 + c_2 + c_3) - (t_s + t_w)(\frac{1}{2}c_1 + \frac{1}{2}c_2 + 2c_3) \\ \hat{t}_{yy}(\mathbf{k}) &= 2t_w(c_1 + c_2 + c_3) - (t_s + t_w)(\frac{3}{2}c_1 + \frac{3}{2}c_2) \\ \hat{t}_{xy}(\mathbf{k}) &= \hat{t}_{yx}(\mathbf{k}) = \frac{\sqrt{3}}{2}(t_s + t_w)(c_1 - c_2) \quad , \end{aligned}$$

\mathbf{k}	θ_1	θ_2	θ_3	c_1	c_2	c_3	$E_+(\mathbf{k})$	$E_-(\mathbf{k})$
Γ	0	0	0	1	1	1	$3(t_s - t_w)$	$3(t_s - t_w)$
K	$\frac{2\pi}{3}$	$\frac{2\pi}{3}$	$-\frac{4\pi}{3}$	$-\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{3}{2}(t_s - t_w)$	$-\frac{3}{2}(t_s - t_w)$
M	π	0	$-\pi$	-1	1	-1	$3t_s + t_w$	$-t_w - 3t_s$

Table 2: Dispersion at high symmetry points.

where $\theta_3 \equiv -(\theta_1 + \theta_2)$ and where $c_j = \cos \theta_j$. Thus,

$$\begin{aligned} \hat{t}_{\mu\nu}(\mathbf{k}) &= \begin{pmatrix} \frac{1}{2}(3t_w - t_s)(c_1 + c_2) - 2t_s c_3 & \frac{\sqrt{3}}{2}(t_s + t_w)(c_1 - c_2) \\ \frac{\sqrt{3}}{2}(t_s + t_w)(c_1 - c_2) & \frac{1}{2}(t_w - 3t_s)(c_1 + c_2) + 2t_w c_3 \end{pmatrix} \\ &= (t_w - t_s)(c_1 + c_2 + c_3) + \frac{1}{2}(t_w + t_s)(c_1 + c_2 - 2c_3) \sigma^z + \frac{\sqrt{3}}{2}(t_s + t_w)(c_1 - c_2) \sigma^x \end{aligned}$$

The eigenvalues of $\hat{H}_{\mu\nu}(\mathbf{k}) = -\hat{t}_{\mu\nu}(\mathbf{k})$ are then

$$E_{\pm}(\mathbf{k}) = (t_s - t_w)(c_1 + c_2 + c_3) \pm (t_s + t_w) \sqrt{c_1^2 + c_2^2 + c_3^2 - c_1 c_2 - c_2 c_3 - c_1 c_3} \quad .$$

Note that under a 60° rotation, $\mathbf{a}_1 \rightarrow -\mathbf{a}_3$, $\mathbf{a}_2 \rightarrow -\mathbf{a}_1$, and $\mathbf{a}_3 \rightarrow -\mathbf{a}_2$, so $(\theta_1, \theta_2, \theta_3) \rightarrow (-\theta_3, -\theta_1, -\theta_2)$. This symmetry is manifestly preserved by the above dispersions.

(b) See the results in Tab. 2. Note that the bands are degenerate at both Γ and K .

(c) See the plot in Fig. 2.

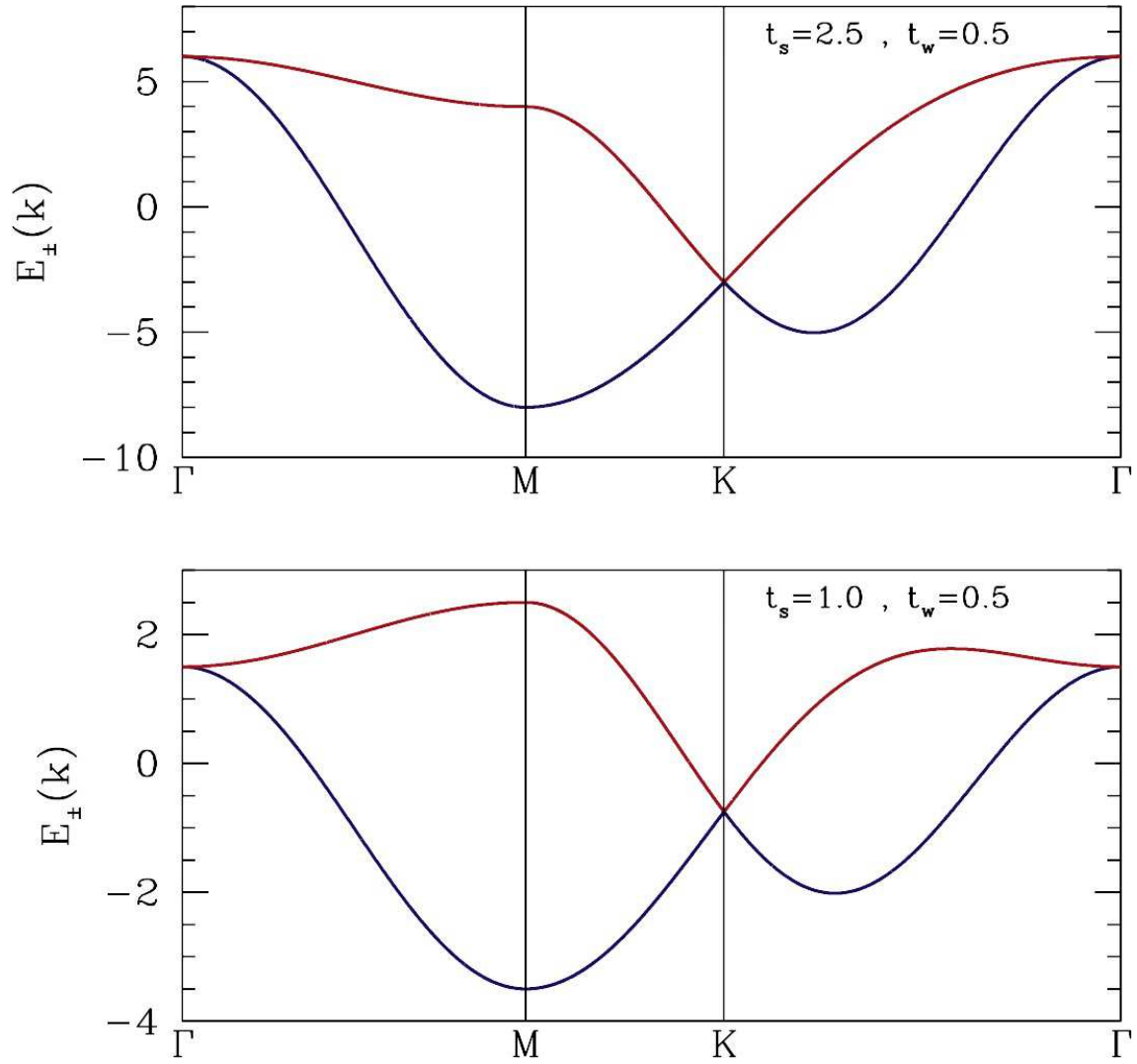


Figure 2: Energy bands $E_{\pm}(\mathbf{k})$ along high symmetry directions for problem 3. Top: $t_s = 2.5$ and $t_w = 0.5$. Bottom: $t_s = 1.0$ and $t_w = 0.5$.