## **Path Length Difference**



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• The interference of two sound waves with identical wavelengths passing through a common point depends on their phase difference there  $\phi$ . If the sound waves were emitted in phase and are traveling in approximately the same direction,  $\phi$  is given by

$$
\phi=\frac{\Delta L}{\lambda}2\pi.
$$

### where  $ΔL$  is their **path length difference**.

• **Fully constructive interference** occurs when  $\phi$  is an integer and multiple of 2π,

$$
\phi = m(2\pi),
$$
 for  $m = 0, 1, 2, ...,$ 

and, equivalently, when ΔL is related to wavelength  $\lambda$  by

$$
\frac{\Delta L}{\lambda} = 0, 1, 2, \dots
$$
 (fully constructive interference).

• **Fully destructive interference** occurs when  $\phi$  is an odd multiple of  $\pi$ ,

$$
\phi = (2m + 1)\pi
$$
, for  $m = 0, 1, 2, ...$ ,

and, equivalently, when ΔL is related to wavelength  $\lambda$  by

$$
\frac{\Delta L}{\lambda} = 0.5, 1.5, 2.5, \ldots
$$
 (fully destructive interference).  
\n
$$
\frac{\Delta L}{\lambda} = 0.5, 1.5, 2.5, \ldots
$$
 (fully destructive interference).



#### **17-4** Intensity And Sound Leve **17-4 Intensity And Sound Level**

• The **intensity** *I* of a sound wave at a surface is the average rate per unit area at which energy is transferred by the wave through or onto the surface

$$
I=\frac{P}{A},
$$

where *P* is the time rate of energy transfer (**power**) of the sound wave and A is the area of the surface intercepting the sound. The intensity *I* is related to the displacement amplitude  $s_m$  of the sound wave by

$$
I=\frac{1}{2}\rho v\omega^2 s_m^2.
$$

#### **17-4** Intensity And Sound Leve **17-4 Intensity And Sound Level**

• The intensity at a distance r from a point source that emits sound waves of power  $P_s$  equally in all directions isotropically i.e. with equal intensity in all directions,

$$
I=\frac{P_s}{4\pi r^2},
$$

where  $4\pi r^2$  is the area of the sphere. 

A point source S emits sound waves uniformly in all directions. The waves pass through an imaginary sphere of radius *r* that is centered on S.



#### **17-4** Intensity And Sound Leve **17-4 Intensity And Sound Level**

### **The Decibel Scale**

The sound level *β* in decibels (dB) is defined as

$$
\beta = (10 \text{ dB}) \log \frac{I}{I_0}.
$$

where  $I_0$  (= 10<sup>-12</sup> W/m<sup>2</sup>) is a reference intensity level to which all intensities are compared. For every factor-of-10 increase in intensity, 10 dB is added to the sound level.

Table 17-2 Some Sound Levels (dB)





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Sound can cause the wall of a drinking glass to oscillate. If the sound produces a standing wave of oscillations and if the intensity of the sound is large enough, the glass will shatter.

#### **17-5** Sources of Musical Sound **17-5** Sources of Musical Sound

Standing sound wave patterns can be set up in pipes (that is, resonance can be set up) if sound of the proper wave- length is introduced in the pipe.

#### **Two Open Ends.**

A pipe open at both ends will resonate at frequencies

$$
f=\frac{v}{\lambda}=\frac{nv}{2L}, \qquad n=1,2,3,\ldots,
$$



#### **One Open End.**

A pipe closed at one end and open at the other will resonate at frequencies



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are nearly equal.

# **17-7** The Doppler Effect

The Doppler effect is a change in the observed frequency of a wave when the source or the detector moves relative to the transmitting medium (such as air). For sound, the observed frequency f' is given in terms of the source frequency f by

 $f' = f \frac{v \pm v_D}{v \pm v_S}$  (general Doppler effect),

where *v* is the speed of sound through the air,  $v<sub>D</sub>$  is the detector's speed relative to the air, and  $v<sub>s</sub>$  is the source's speed relative to the air. In the numerator, the plus sign applies when the detector moves toward the source and the minus sign applies when the detector moves away from the source.

In the denominator the minus sign is used when the source moves toward the detector, the plus sign applies when the source moves away from the detector. 

$$
3\sqrt{1+1}t
$$
  
\n
$$
3\sqrt
$$



**17-7** The Doppler Effect

## **Detector Moving Source Stationary**

## **Source Moving Detector Stationary**



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#### **17-8** Supersonic Speeds, Shock Waves **17-8** Supersonic Speeds, Shock Waves

If the speed of a source relative to the medium exceeds the speed of sound in the medium, the Doppler equation no longer applies. In such a case, shock waves result. The half-angle  $\vartheta$  of the Mach cone is given by

$$
\sin \theta = \frac{vt}{v_S t} = \frac{v}{v_S}
$$
 (Mach cone angle).

A source *S* moves at speed  $v_s$  faster than the speed of sound and thus faster than the wavefronts. When the source was at position  $S_1$  it generated wavefront  $W_1$ , and at position  $S_6$  it generated  $W_6$ . All the spherical wavefronts expand at the speed of sound *v* and bunch along the surface of a cone called the Mach cone, forming a shock wave. The surface of the cone has half-angle  $\vartheta$  and is tangent to all the wavefronts. 



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Speed of sound on air  $V = \sqrt{\frac{B}{D}}$   $B = -V\frac{dV}{dV}$  $P = nRT$   $-\frac{dP}{dV} = P -2$ <br>U= $\sqrt{\frac{P}{S}}$   $-\frac{dP}{dV} = \frac{V}{T} = 1.2$  $P_{air} = 1.225 kg/m^3$ <br>  $P = 1.01x105 Pa$  $V = \sqrt{\frac{L225}{10 \times 10^5}}$  m/s = 287 m/s - 340<u>m</u>  $PV^{\tau}$ = cist =>  $P = \frac{2ms^{\tau}}{\sqrt{s}}$  $B = -VdP$ <br> $V_s = VdP$ <br> $V_s = Vg.287m/s$  $r = \frac{C_0}{C_0} = \frac{C_{v+R}}{C_0} = 1 + \frac{R}{C_0} = 1 + \frac{R}{2R} = 1 + \frac{2}{5} = \frac{2}{5}$ 

$$
V = \sqrt{\frac{2P}{S}}
$$
\n
$$
P = \frac{nRT}{V} \qquad S = \frac{N}{V} = \frac{nNAm}{V}
$$
\n
$$
P = \frac{Q}{Nam} \qquad P = \frac{S}{Nam} \qquad h_0NAT \Rightarrow
$$
\n
$$
P = \frac{Q}{Nam} \qquad R = \sqrt{\frac{S}{Nam}} \qquad h_0NAT \Rightarrow
$$

## **17** Summary

### Sound Waves

Speed of sound waves in a medium having bulk modulus and density

$$
v = \sqrt{\frac{B}{\rho}}
$$
 Eq. (17-3)

### Interference

If the sound waves were emitted in phase and are traveling in approximately the same direction,  $\phi$  is given by

$$
\phi = \frac{\Delta L}{\lambda} 2\pi, \qquad \text{Eq. (17-21)}
$$

## Sound Intensity

The intensity at a distance r from a point source that emits sound waves of power Ps is

Eq. (17-3) 
$$
I = \frac{P_s}{4\pi r^2}.
$$
 Eq. (17-28)

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### Sound Level in Decibel

The sound level b in decibels (dB) is defined as  $=$  (10 dB)  $\log \frac{1}{2}$  (17-2)  $\left(1\right)$   $\left(1\right)$   $\left(1\right)$   $\left(1\right)$   $\left(1\right)$ where  $I_0$  (=  $10^{-12}$  W/m<sup>2</sup>) is a reference intensity **Eq. (17-29)**

## **17** Summary

### **Standing Waves in Pipes**

A pipe open at both ends

$$
f = \frac{v}{\lambda} = \frac{nv}{2L}, \qquad n = 1, 2, 3, \dots, \text{ Eq. (17-39)}
$$

A pipe closed at one end and open at the other

$$
f = \frac{v}{\lambda} = \frac{nv}{4L}, \quad n = 1, 3, 5, \dots.
$$
 Eq. (17-41)

### The Doppler Effect

• For sound the observed frequency  $f'$  is given in terms of the source frequency f by

$$
f' = f \frac{v \pm v_D}{v \pm v_S}
$$
 Eq. (17-47)

### l The half-angle θ of the Mach cone is given by

Sound Intensity

$$
\sin \theta = \frac{v}{v_S} \qquad \text{Eq. (17-57)}
$$

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In Maxwell's time (the mid 1800s), the visible, infrared, and ultraviolet forms of light were the only electromagnetic waves known. Spurred on by Maxwell's work, however, Heinrich Hertz discovered what we now call radio waves and verified that they move through the laboratory at the same speed as visible light, indicating that they have the same basic nature as visible light. As the figure shows, we now know a wide spectrum (or range) of electromagnetic waves: Maxwell's rainbow.<br>
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